

Re-stating the PBL parameters in off-line dispersion models

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Motivation: off-line dispersion modelling

- No direct connection between dispersion and NWP models
- pro-s
 - multiple dispersion runs are cheap (no need to rerun NWP)
 - possibility for multi-NWP-input
 - easier model development (smaller code, simpler portability issues, etc)
 - no influence of simplifications in one model to another one
- contra-s
 - no possibility for feedback to NWP
 - internal NWP variables are not accessible
 - limitations on input data size force compromises on accuracy and level of details
- Necessity for special module: meteorological preprocessor



Motivation: meteorological pre-processor

- Prepares meteorological input data for dispersion model
 - extra variables, non-existing in the input files
 - checking/restating the governing equations as they are in the dispersion model
 - enhanced resolution in time and/or space
- Varying levels of complexity
 - simple interpolation & range-checking
 - sophisticated algorithms up to own assimilation of meteorological observations and recomputation of dynamic equations (MM5)



Motivation: boundary layer parameters

- Numerous approaches to parameterization
- Specific variables and equations vary from model to model and even from run to run
- Most of ABL parameters are not explicitly validated in NWP models and not available in the output files
- Result: practically all dispersion models include re-stating the ABL basic parameters in their meteorological pre-processor



Problem statement

- Available: profiles of basic meteorological variables: wind \vec{u} , temperature *T*, humidity *q*
- Find: basic ABL parameters: temperature, velocity and humidity scales *T*, *u**, *q**, Monin-Obukhov length *L*, profile of some characteristic of turbulence, e.g. *K_Z* – if K-theory is used
- Verification possibility: consistency checking via comparison of sensible and latent heat fluxes $H_{\rm S}$, $H_{\rm I}$



Problem solution(1)

• Closure equation obtained from M-O similarity consideration (Berlyand & Genikhovich, 1971):

$$\sqrt{TKE} \, \frac{d}{dz} \left(\frac{K_z}{TKE} \right) = \kappa \, c_{\varepsilon}^{0.25} \Phi(\zeta)$$

- For practical applications, $\Phi = 1$; it corresponds to a differential expression combining the eddy diffusivity and TKE
- Using this closure expression together with equations governing the surface layer, one can obtain the following formula:

$$K_{z} = \left\{ \frac{\kappa}{2} \int_{0}^{z} \frac{\left[\left(\frac{dU}{dz} \right)^{2} - \sigma \beta \frac{d\theta}{dz} \right]^{5/4}}{\left(\frac{dU}{dz} \right)^{2} - 0.5\sigma \beta \frac{d\theta}{dz} dz} dz \right\}^{2}$$



Problem solution (2)



Here all derivatives are NOT computed numerically but rather taken from the analytical approximations of profiles.

Since $z_k \sim 1m$, these profiles can be taken purely logarithmic. Non-logarithmic corrections start to play a strong role at $|z/L| \sim 0.5$

Assuming the logarithmic shape, it is enough to have 2 values – at the screening and the 1st model levels – to determine the profile.

All fluctuating and not well-defined parameters are inside the integral, thus their effect is smoothed out



Problem solution (3)

Logarithmic profile assumption near the surface (z < < |L|)

$$T(z) = T_g + T^* \operatorname{Pr} \ln \frac{z}{z_{0T}} \implies \frac{\partial T}{\partial z}(z) = T^* \operatorname{Pr} \frac{1}{z}$$

Having temperature values at two levels, obtain:

$$T(z_2) - T(z_1) = T^* \operatorname{Pr} \ln \frac{z_2}{z_1} \implies T^* \operatorname{Pr} = \frac{T(z_2) - T(z_1)}{\ln \frac{z_2}{z_1}} \implies \frac{\partial T}{\partial z} = \frac{T(z_2) - T(z_1)}{z \ln \frac{z_2}{z_1}}$$

Analogously, for velocity scale:

 $\frac{\partial U}{\partial z} = \frac{U(z_2) - U(z_1)}{z \ln \frac{z_2}{z_1}}$

These dependencies are substituted into K_z formula, where the integral is tabulated



Method verification: measurements

Eddy-correlation measurements, Tsimlyansk, 1976



Groisman & Genikhovich (1997), using the lower available measurement level and ground surface; the temperature jump is estimated after Zilitinkevich (1970)



Profile measurements, Cabauw, 1987





Problem discussion: iterative solution



- differentiations have to be done numerically
- convergence of the iterations is not proven



Comparison of solutions



Friction velocity, iterative solution

Friction velocity, K_z-based solution



Problem discussion (2)

 A supposedly simpler way is to use NWP fluxes of momentum (M_x, M_y) and heat (H_s, H_l) to get the ABL characteristics, for example:

$$u^{*} = \left[\frac{1}{\rho}\sqrt{M_{x}^{2} + M_{y}^{2}}\right]^{1/2}$$
$$T^{*} = \frac{-H_{s}}{c_{p}\rho u^{*}} \qquad q^{*} = \frac{-H_{l}}{E\rho u^{*}}$$

 $L = \frac{T_{sl} (u^*)^2}{\kappa \, gT^*}$

Here ρ is an air density, c_{ρ} is a constant-pressure heat content, *E* is an energy of evaporation or sublimation, depending on screen-level (usually, 2m) temperature T_{sl} , *k* is a von-Karman constant, *g* is an acceleration of gravity

A problem: the heat and momentum fluxes are not routinely verified in the NWP models and thus cannot be considered reliable



Comparison with NWP (HIRLAM, ECMWF)

- Intuitively, there must be almost 1:1 agreement
 - theoretical basis is more or less the same, variations in the formulations should not lead to excessive quantitative discrepancies
 - within HIRLAM *u,q,T* profiles and heat fluxes are computed together, thus being highly correlated
- However, certain deficiencies are inevitable
 - still, there are differences in the computational algorithms
 - HIRLAM & ECMWF provide accumulated fluxes e.g. for 3 hours, while *u,q,T* are instant, thus re-stated fluxes will be instant too



Comparison with HIRLAM





Verification statistics: HIRLAM, Jan-March 2000, night





Verification statistics: HIRLAM, May-Sep 2000, day





Verification statistics: time correlation, quantile charts





Verification statistics: stratification influence

Ratio SILAM / NWP for sensible heat flux, South Atlantic



A ratio of re-stated and NWP-computed sensible heat fluxes versus re-stated one. HIRLAM 5.2.1, south Atlantic, complete year 2000



Comparison of time series (latent flux)





Comparison of time series (sensible flux)





Discussion of comparison

- The closest values are shown for latent heat flux; usually reasonable similarity of sensible heat fluxes over sea and in unstable cases (with a correction of *Pr*-number dependence on stratification)
- Larger deviations in stable cases
 - slightly stable case is practically OK
 - stronger stability leads to re-stated downward flux much stronger than HIRLAM one
- Fluxes should not be the same (above reasons)
- High sensitivity of the parameters "by nature"
 - small differences in approaches may magnify
 - possibility for explosion of numerical error (hardly)
 - coding error in implementation in either model





An effect of "error explosion" for ECMWF operational model data, 1-17.12.2000



Call for future studies

- Available methodology
 - universal approach for re-stating the main ABL characteristics from the basic meteorological variables
 - verification against observations showed good results
- Existing problem
 - comparison with HIRLAM heat fluxes showed significant differences, especially for sensible heat flux over terrestrial areas in stable cases. Reason is largely unknown
- Research needed
 - comparison with independent datasets
 - ECMWF model fields
 - Sodankyla mast data
 - other datasets
 - fine-tuning of the methodology (and/or HIRLAM) and/or its implementation