<u>Atmospheric Planetary Boundary</u> Layers (ABLs / PBLs) in stable, neural and unstable stratification: scaling, data, analytical models and surface-flux algorithms

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Part II: The role of large, organised eddied in the convective heat/mass transfer

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Content

Large eddies (convective winds) enhance mixing Homogeneous surfaces:

- Shear-free convection \rightarrow cells
- Sheared convection \rightarrow rolls

Heterogeneous surfaces and heat islands:

- polynias / big cities
- leads

<u>Future work:</u> improved CBL module and fluxparameterization









Closed cloud cells over the Atlantic Ocean



Low wind convection

Radiative cooling of the upper boundary of clouds causes narrow cold descending plumes surrounded by warmer updraughts









Open cloud cells over the Pacific Ocean



Low wind convection

Narrow warm uprising plumes surrounded by colder downdraughts, driven by the positive buoyancy flux over warm sea surface









Cloud "streets" over the Amazon River <u>Strong wind: shear-generated convective rolls</u>



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Convective-shear waves



Perturbation analysis predicts generation of roll-type structures in the plane perpendicular to the mean wind (Elperin et al.)









Turbulent convection in laboratory

Semi-organised circulation in a box with heated bottom (Rayleigh-Benard apparatus): vertical (left) and horizontal (right) cross-sections (Elperin et al.)



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SHEAR-FREE CONVECTION OVER WIDE HEATED AREAS: THE INFLUENCE OF SEMI-ORGANISED EDDIES ON THE HEAT / MASS TRANSFER AT THE SURFACE

Large buoyancy-driven eddies in CBLs ↓
"Convective winds": convergence IBL-type flows towards plume axes near the surface ↓ Local sears in IBLs enhance mixing ↓ Stronger turbulent fluxes







Large convective eddy



Williams and Hacker (1992) airborne measurements: Arrows show the large-eddy velocity field (subtracting mean wind). Solid lines show deviations of potential temperature θ from its large-eddy averaged value < θ >. The iso-surface θ -< θ >=0 marks the side walls of the updraught.



MPRAS **A** 1



LES "portrait" of the 1st characteristic (most energetic) eddy in shear-free CBL



Solid contours in horizontal plane mark up-draughts with w from 0.14 to 0.84 m/s. The maximum updraught velocity is 1.2 m/s, maximum horizontal velocity is 2.8 m/s. Dashed contours mark downdraughts (varying from -0.16 to -0.8 m s-1). Bold curves with arrows in y,z plain show streamlines. The domain size is given in km.

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Horizontal cross-sections of plumes and downdraughts over smooth and rough surfaces upper snapshots – surface layer, below – upper CBL; left – smooth, right – very rough



Prior large-eddy shear models

- Deardorff's scaling: CBL depth *h*, convective velocity $W_* = (\beta F_{\theta sx} h)^{1/3}$
- Resistance coefficients: U_*/W_* , $F_{\theta sx}/(W_* \Delta \theta)$ versus h/z_{0u}
- Schumann (1988) 1-layer model, dominant role of buoyancy forces
- Sykes et al. (1993) 1-layer model, dominant role of large-eddy shears
- Zilitinkevich, Grachev & Hunt (1998) 2-layer buoyancy + shear model
- Limited applicability; insufficient accuracy over very rough surfaces







New model

- Large eddy life time considerably exceeds $h/W_* \sim 20$ min Large-scale convergence flows are internal boundary layers (IBLs): $W_*dh_I/dx=0.24W_c(z=h_I)$
- Locally equilibrium turbulence determines *x*-dependent values of $\tau(z) = u_*^2(x), F_{\theta sx}(x), \text{ large-eddy MO length: } L(x) = \tau^{3/2} (\beta F_{\theta sx} h)^{-1}$ Eddy viscosity / conductivity $K \sim z[u_* + Constant W_*(z/h)^{1/3}]$ Solution for area-averaged fluxes: $U_*^4 = \langle \tau^2 \rangle, F_{\theta s} = \langle F_{\theta sx} \rangle$ <u>Very rough surfaces:</u> surface shear layer diminishes effective roughness length: $z_0/z_{0u} = 1 + C_{0C}(z_{0u}/L)^{1/3}$







IBL within the large eddy: (a) low roughness

Typical height of roughness elements smaller than the MO length scale Two-layer structure: logarithmic + free-convection vertical profiles



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IBL within the large eddy: (b) large roughness

Typical height of roughness elements larger than the MO length scale One-layer structure; essential stability dependence of the roughness length



Resistance coefficient U_*/W_* vs. h/z_{0u}

Blue symbols show field data; red symbols - LES (NERSC data with error bars)



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Heat transfer coefficient U_{*}/W_{*} vs. h/z_{0u}

Blue symbols show field data; red symbols - LES (NERSC data with error bars)



Comparison of heat transfer models

Solid line – our model validated against field data and LES (CBLs over natural rough surfaces); doted line – classical heat transfer law: Nu=0.14 Ra^{1/3} (lab experiments)



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Conclusions and future work

Failure of the classical local theory

Key role of organised structures:

- not to be confused with largest turbulent eddies
- notice inverse energy cascade

New physics at large roughness:

- Diminishing logarithmic sub-layer
- <u>Stability dependence of the effective roughness length</u>

Future work

- Theory and its validation for sheared convection
- Improved parameterization of surface fluxes







CHAOS GENERATES ORDER

Large-scale convective structures are not turbulence: they are essentially regular motions fed by the convective energy production through the inverse energy cascade

over heat island

over vide heated area









